

# Towards Distribution Amplitude and GPD Calculations on a Quantum Computer

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Parton Distributions---  
a Great Challenge Both in  
Exp. and Th.

LaMET (Ji, 2013) is the most popular  
theoretical approach so far.

# This year's highlights

- Precision frontier: Coulomb-gauge, Renormalon/threshold resummation/long-distance extrapolation/higher loop matchings, enhanced Interpolators...
- Higher dimensional frontier: TMD/CS kernel (continuum/physical-point), GPD non-zero skewness,  $\Lambda$ -Baryon DA...
- Flavor frontier: gluonic helicity/PDF/CS, heavy flavor DAs...

# Lattice QCD

- Feynman's Path Integral, Monte Carlo, sign problem and real time dependence challenging

# An Alternative: Quantum Computation

- Studying how a state evolves under unitary transformation, including real time evolution
- Matrix elements of non-unitarity operators embedded into unitary transformations of larger (Hilbert) spaces with ancillary qubits added.
- Dealing with real time evolution and systems with the sign problem in classical computation in the same way with other problems

# Why Quantum Computation?

- **Memory:**

**Classical:** n bits can store an integer from 0 to  $2^n - 1$ . **Quantum:** n-qubits can store  $2^n - 1$  complex numbers.

- However, if **entanglement** is lost (due to **decoherence**) such that the n-qubit state becomes a cross product state, then only n-1 complex numbers can be stored.
- **Superposition** allows the  $2^n$  states/**solutions** to be studied at the same time.

# Renormalon Ambiguity in LaMET

$$\tilde{q}(x, P_z) = \int_{-1}^1 \frac{dy}{|y|} C\left(\frac{x}{y}, \frac{\mu}{yP_z}\right) q(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{(xP_z)^2}, \frac{\Lambda_{\text{QCD}}^2}{((1-x)P_z)^2}\right)$$

- OPE: separation of long and short distance physics not strict in MS-bar scheme
- Matching kernel: IR Borel resummation leads to  $\Lambda_{\text{QCD}}^2/P_z^2$  ambiguity to be canceled by power corrections
- Large momentum expansion fails when the parton momentum is small

# Quantum PDF is IR Renormalon Ambiguity Free

- No IR renormalon ambiguity. Can access  $x$  near 0 and 1.
- No IR renormalon in GPD. What about TMD?

# Other Possible Advantages of Quantum PDFs

- Accessing nonvalance partons with same expanses as valence ones (no disconnected diagram complications as in lattice QCD)
- 3+1 d QCD at end points of  $x$ --- a clear demonstration of quantum advantage on a scientifically important problem

# Toy: PDF in QED<sub>2</sub>

- Linear, confining potential. Positronium mimics meson in 3+1D QCD
- Working in Minkowski space, accessing lightcone correlators directly

$$\mathcal{O}(z_1, z_2) = \bar{\psi}(z_1 n) n \cdot \gamma W_n(z_1 n, z_2 n) \psi(z_2 n). \quad n^\mu = (1, -1)$$

$$D(z_1 - z_2) = \frac{1}{2n \cdot P} \langle P | \mathcal{O}(z_1, z_2) | P \rangle$$

$$f(x) = \int_{-\infty}^{\infty} \frac{dz n \cdot P}{2\pi} e^{-izn \cdot Px} D(z)$$

Super renormalizable

$$f(x, a) = f(x) + \mathcal{O}(a)$$

# Kogut Susskind Fermions in 1+1 d

$$\psi(x) = (\phi(n), \phi(n+1))^T$$

$$\{\phi(n), \phi^\dagger(k)\} = \delta_{n,k} \quad , \quad \{\phi(n), \phi(k)\} = 0 \quad .$$

(similar to 3+1d equal time  $\{\psi_a(\mathbf{x}), \psi_b^\dagger(\mathbf{y})\} = \delta^{(3)}(\mathbf{x} - \mathbf{y}) \delta_{ab}$ .)

$$\hat{H} = \frac{i}{2a_n} \sum_n \left[ \phi^\dagger(n) \phi(n+1) - \phi^\dagger(n+1) \phi(n) \right]$$

$$i \frac{d}{dt} \phi(n) = \left[ \phi(n), \hat{H} \right] = \frac{i}{2a_n} ( \phi(n+1) - \phi(n-1) )$$

$$\frac{d}{dt} \psi(x) = \alpha_x \frac{d}{dx} \psi(x) \quad , \quad \alpha_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\gamma_0 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \quad , \quad \gamma_x = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad , \quad \{\gamma_\mu, \gamma_\nu\} = 2g_{\mu\nu}$$

# Jordan Wigner Mapping

Occupied (unoccupied) fermion state at site  $n$  as spin up (down)

$$\phi(n) = \left(-i\hat{Z}_1\right) \left(-i\hat{Z}_2\right) \dots \left(-i\hat{Z}_{n-1}\right) \sigma_n^-$$

$$\phi^\dagger(n) = \left(+i\hat{Z}_1\right) \left(+i\hat{Z}_2\right) \dots \left(+i\hat{Z}_{n-1}\right) \sigma_n^+$$

$$\phi(0) = \sigma^- \otimes I \otimes I \otimes I, \quad \phi^\dagger(0) = \sigma^+ \otimes I \otimes I \otimes I$$

$$\phi(1) = -i\hat{Z} \otimes \sigma^- \otimes I \otimes I, \quad \phi^\dagger(1) = +i\hat{Z} \otimes \sigma^+ \otimes I \otimes I$$

$$\phi(2) = -\hat{Z} \otimes \hat{Z} \otimes \sigma^- \otimes I, \quad \phi^\dagger(2) = -\hat{Z} \otimes \hat{Z} \otimes \sigma^+ \otimes I$$

$$\phi(3) = +i\hat{Z} \otimes \hat{Z} \otimes \hat{Z} \otimes \sigma^-, \quad \phi^\dagger(3) = -i\hat{Z} \otimes \hat{Z} \otimes \hat{Z} \otimes \sigma^+$$

$$\{\phi(1), \phi^\dagger(1)\} = I \quad \{\phi(1), \phi(1)\} = 0 \quad \{\phi(0), \phi^\dagger(1)\} = 0$$

$$\hat{H} = -\frac{1}{2a} \sum_n \left[ \sigma_n^+ \sigma_{n+1}^- + \sigma_{n+1}^+ \sigma_n^- \right] + \frac{m}{2} \sum_n (-)^n \hat{Z}_n$$

# Hamiltonian of QED<sub>2</sub>

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}(i\not{D} - m)\psi, \quad D_\mu = \partial_\mu + ieA_\mu$$

$$\Pi^\mu = \frac{\partial\mathcal{L}}{\partial(\partial_0 A_\mu)} = F^{\mu 0}, \quad \Pi^0 = 0 \text{ implies } A^0 \text{ is not dynamical.}$$

$$\frac{\partial\mathcal{L}}{\partial A_0} = J^0 - \partial_\mu \Pi^\mu = 0 \quad (\text{Gauss's law, a quantum constraint})$$

$A^1=0$  (the axial gauge) (open boundary condition in x)

$$\begin{aligned} \mathcal{H} &= \pi\partial_0\psi - \mathcal{L} \\ &= \frac{1}{2}E^2 - i\bar{\psi}\gamma^1\partial_1\psi + m\bar{\psi}\psi, \end{aligned}$$

# Wilson Line

$$W_n(z_1 n, z_2 n) = e^{-ie \int_{z_2}^{z_1} n \cdot A(z) dz}$$

Gauge invariant

$$O'(x, y) = \bar{\psi}(x) n \cdot \gamma Q_n(x) \bar{Q}_n(y) \psi(y)$$

$$\mathcal{L}' = \mathcal{L} + \bar{Q}_n(x) i n \cdot D Q_n(x)$$

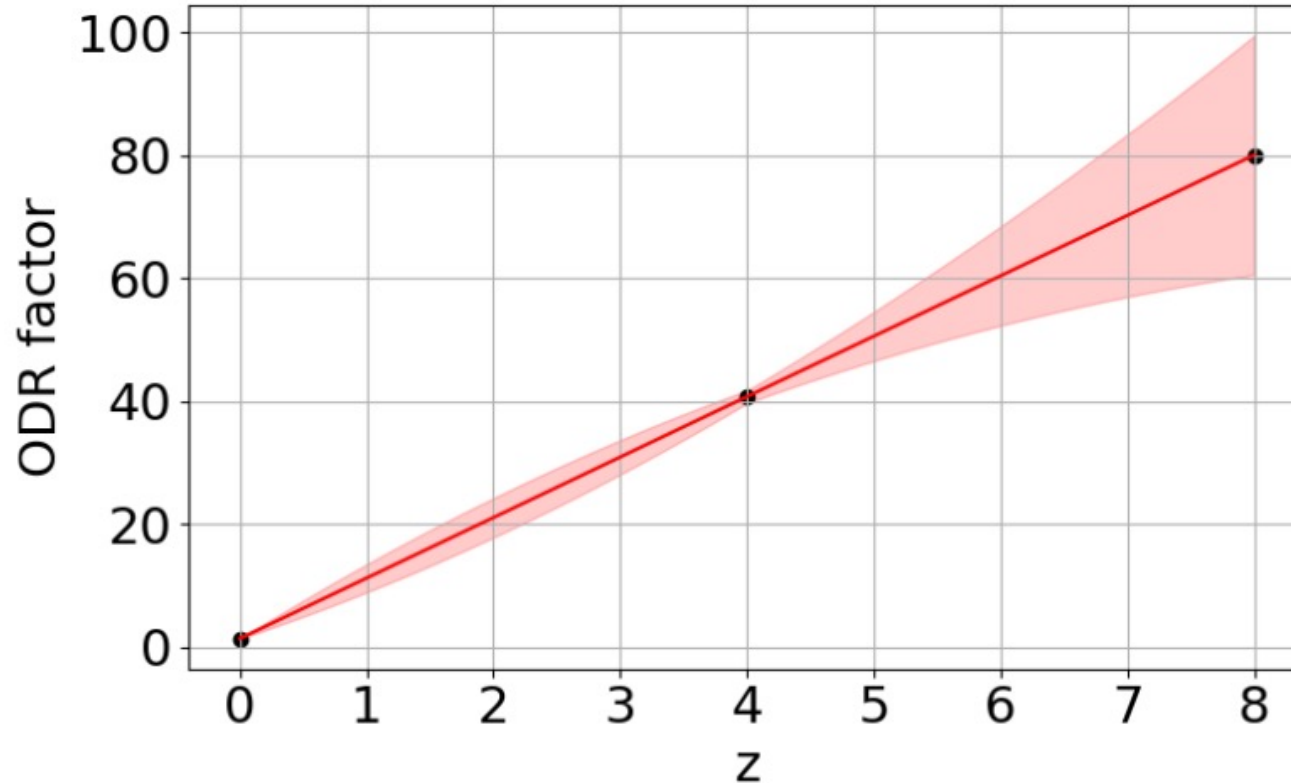
$$\delta J^\mu = -e \bar{Q}_n(x) n^\mu Q_n(x)$$

Wilson line carries the same charge as the fermion

# Quantum Computation

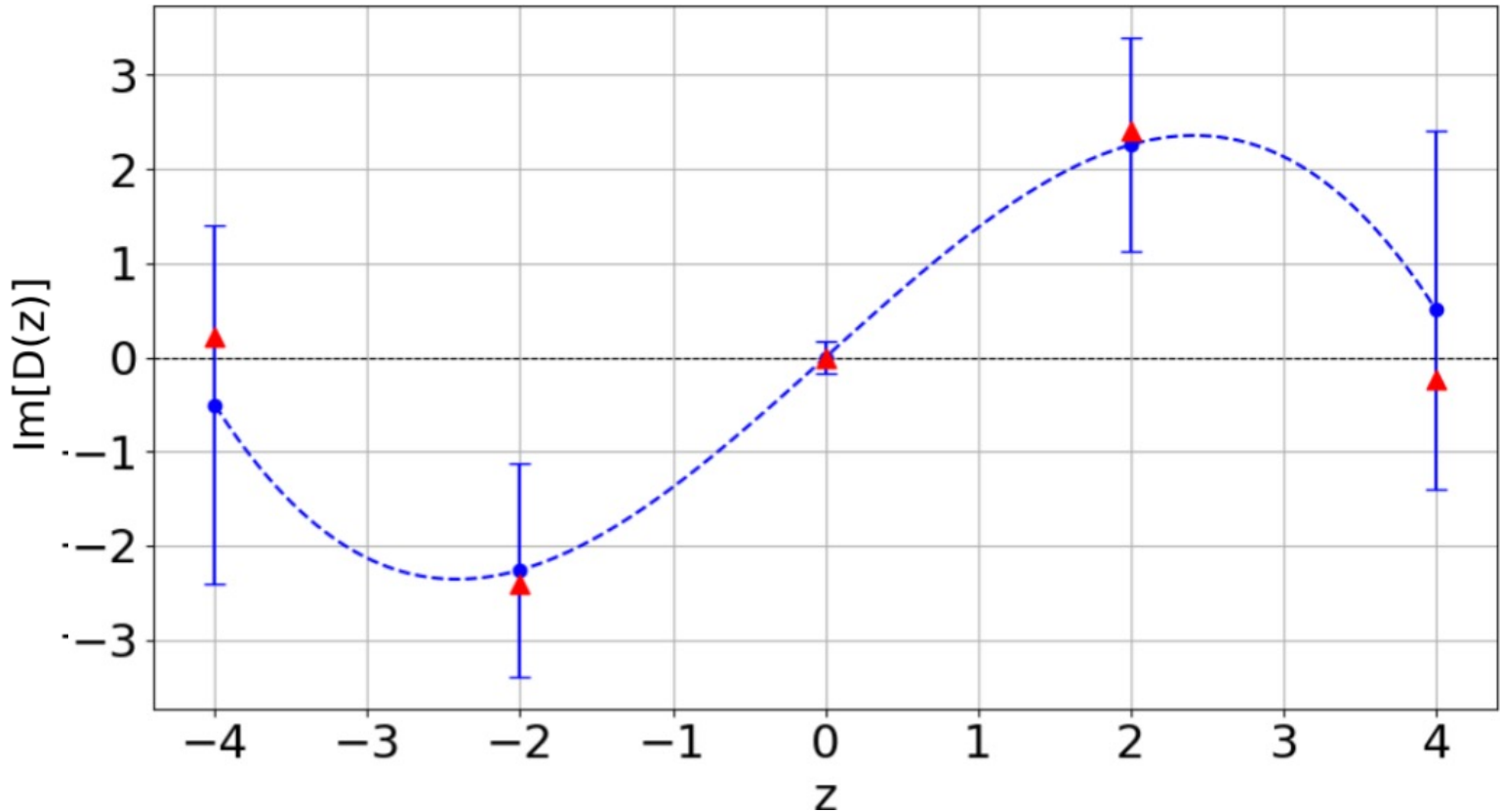
- IBM cloud quantum computers
- 11 qubits: 5 physical spatial sites, 10 staggered fermion sites, 1 ancillary qubit; Zigzag Wilson lines
- 2-qubit gate depth  $\sim 5K$  reduced to  $\sim 500$  (very critical!). Mainly from changing state preparation, operator choice, and doubling the step size for time evolution and Wilson line. Each reduces the gate depth by half.

# Error Mitigations



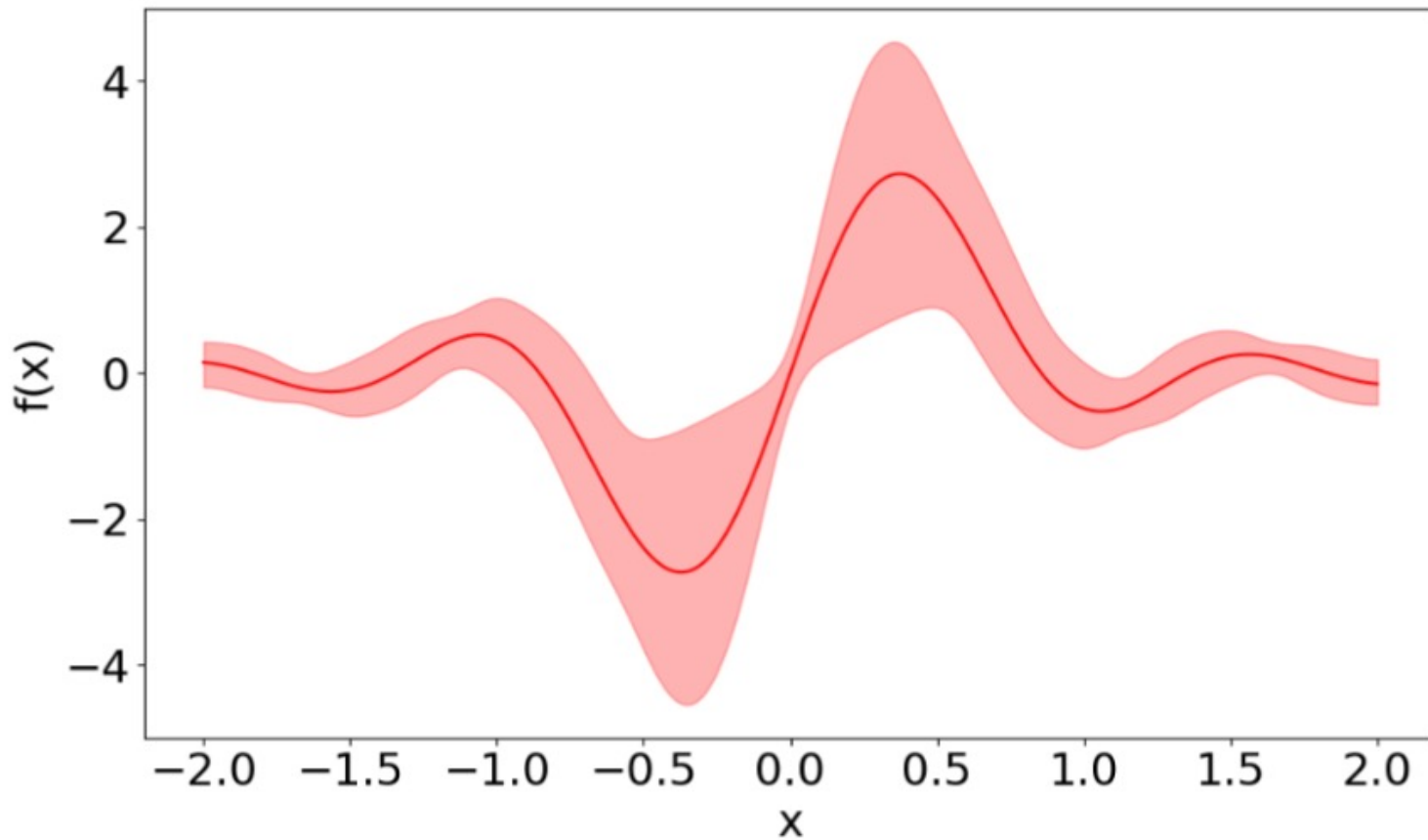
Dynamical decoupling (DD), Pauli twirling (PT) with randomization = 100, and Operator Decoherence Renormalization (ODR).

# Result: Lightcone Correlator



8,000 shot result from IBM quantum computers.  
Triangles from a classical simulator.

# First “Quantum” PDF Result



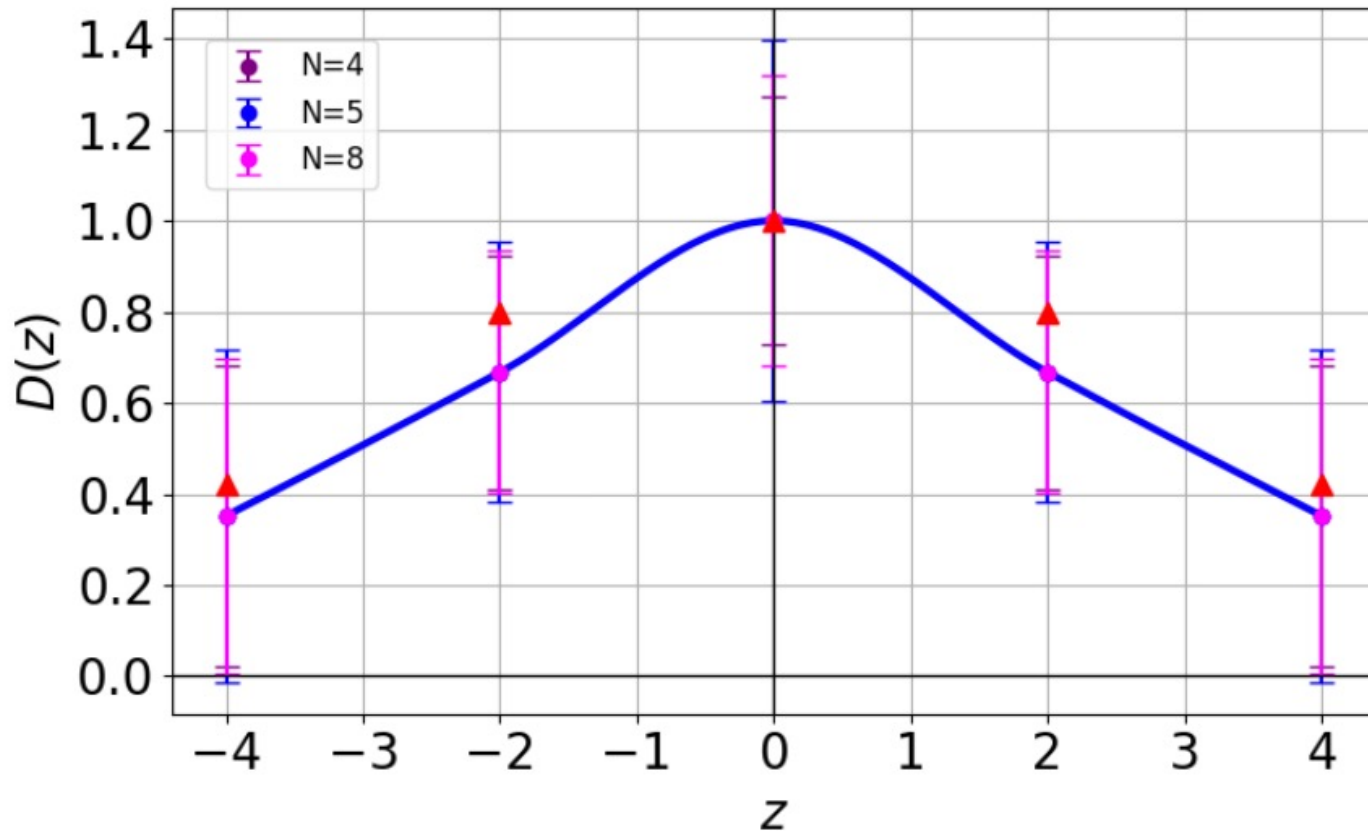
$$\int_{-1}^1 dx x f(x) = 1.04^{+0.66}_{-0.76},$$

$$\int_0^1 dx f(x) = 1.37^{+0.89}_{-1.02}.$$

# Scaling

- 2 qubit gate number scales as the cubic of the volume. This calculation is state of the art already.
- (1+1) d QED can be completely solved with machines available in 2029.

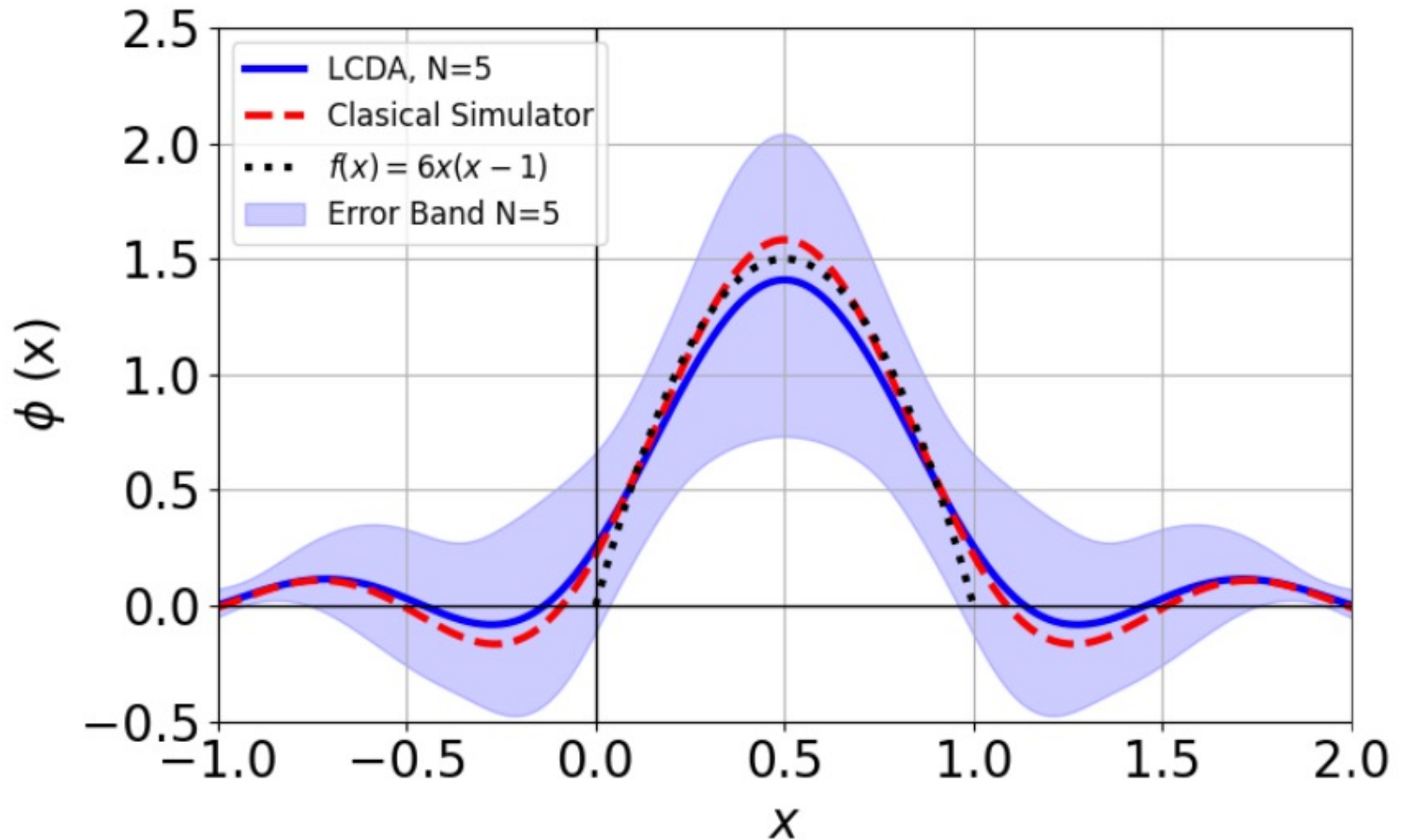
# Meson DA



$\langle \Omega | \mathcal{O} | P \rangle$

Two qubit gate/depth  $\sim 1,000$

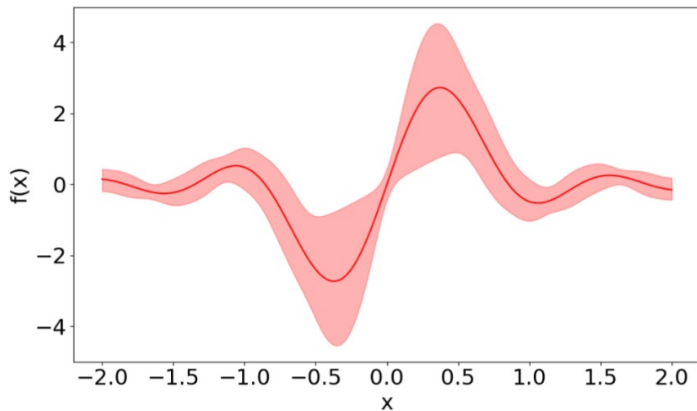
# First “Quantum” DA Result



# GPD?

$$H(x, \xi, t) = \frac{1}{2P^+} \int \frac{dz^-}{2\pi} e^{ixP^+z^-} \langle M(q) | \bar{\psi} \left( -\frac{z^-}{2} \right) \gamma^+ \mathcal{W} \psi \left( \frac{z^-}{2} \right) | M(0) \rangle$$

$$t = -\frac{4M^2\xi^2}{1-\xi^2}.$$



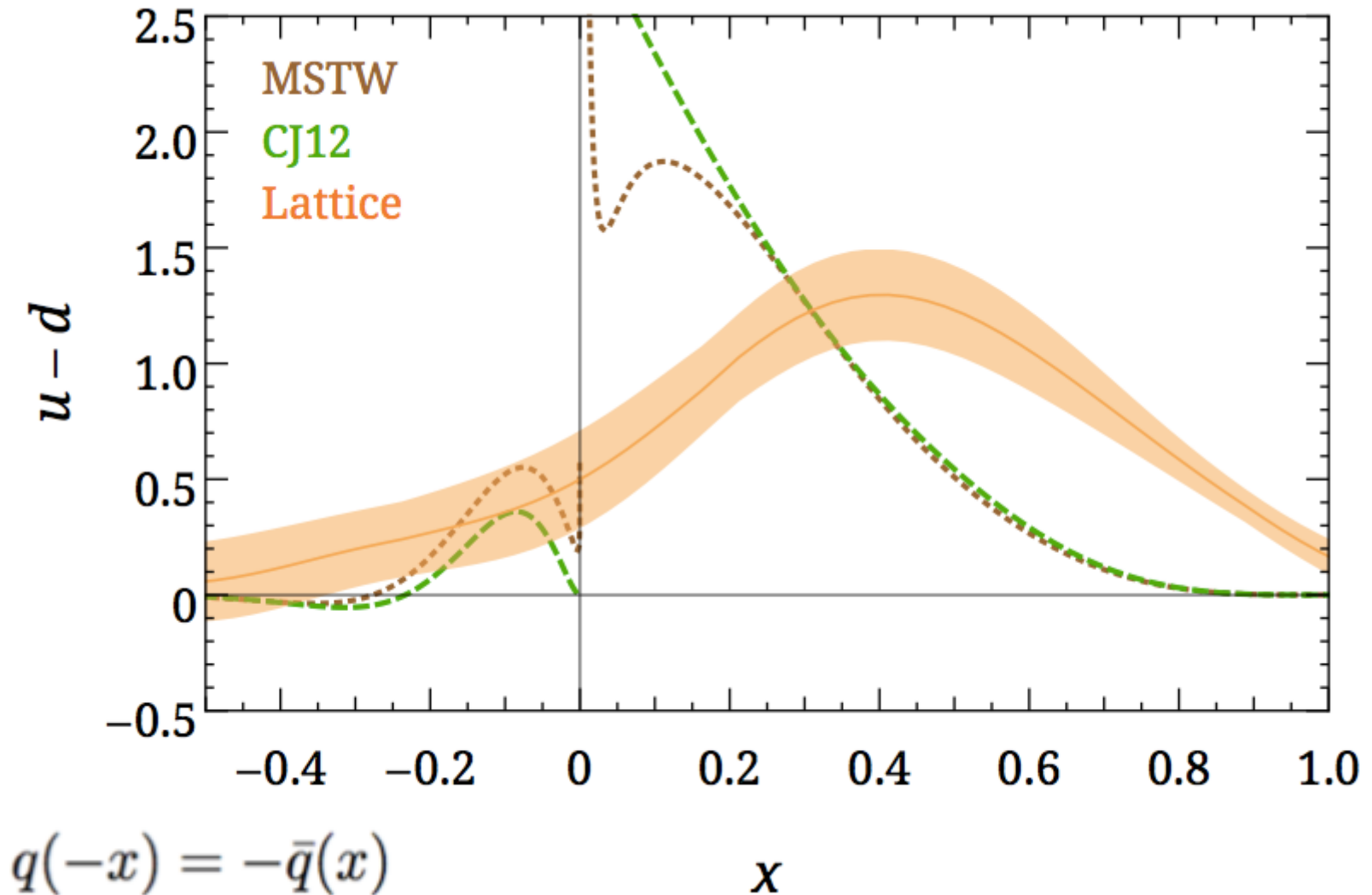
Two qubit gate/depth  $\sim 8,000$

Cannot do it **now**.

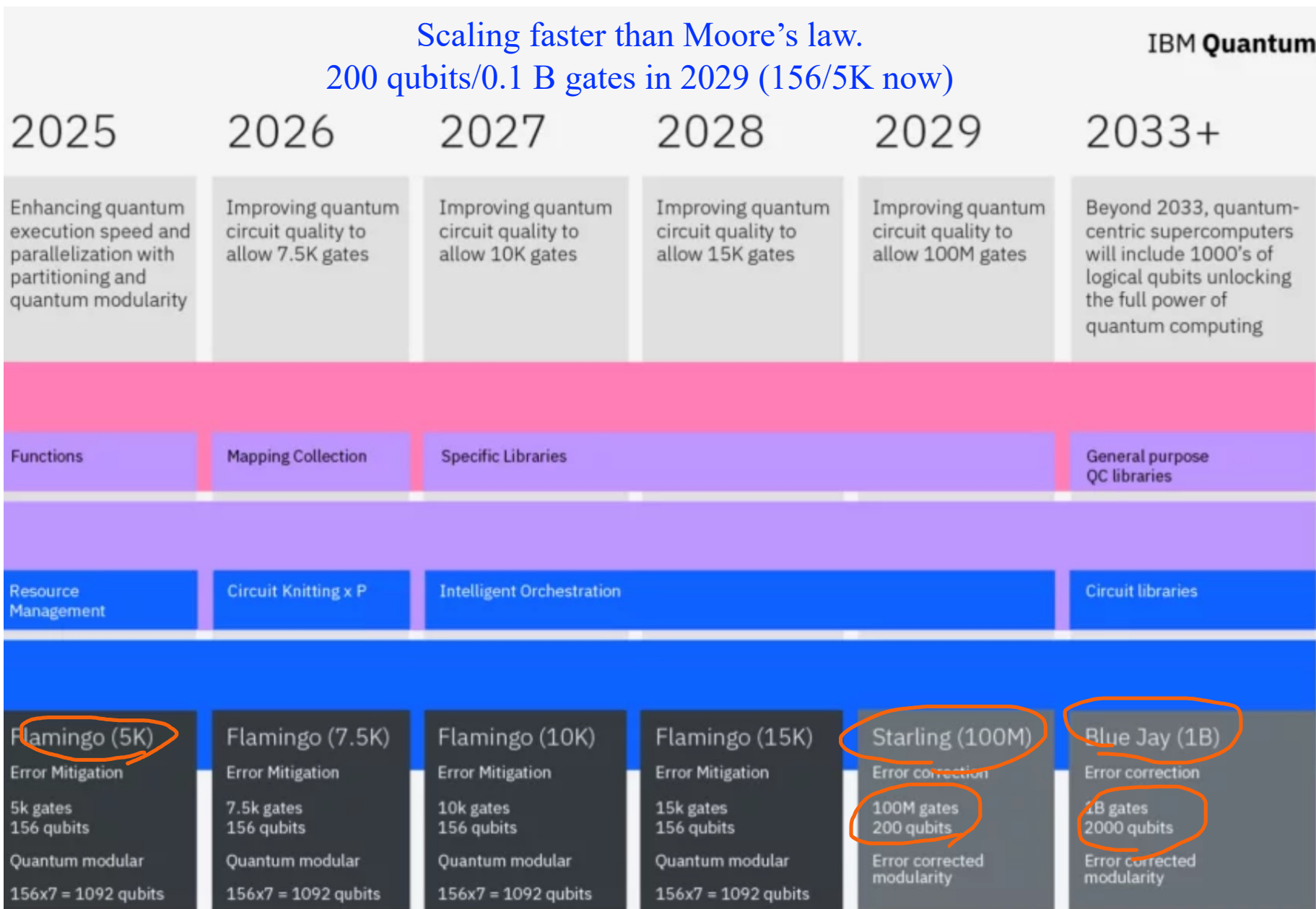
# Outlook

# 1<sup>st</sup> LaMET Calculation

Lin, JWC, Cohen, Ji (1402.1462)



# IBM Quantum Roadmap



See what will happen in 12 years!

# Backup

# Remaining Errors

- Statistical errors (standard errors): inversely proportional to the number of shots
- Remove finite volume effect and lattice spacing, need  $\sim 3$  and  $\sim 2$  times more sites
- 2-qubit gate depth scales with 3 powers of the spatial volume
- Machines available in 2029!